

Monte Carlo Simulation of Positronium Cooling for Bose-Einstein Condensation

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Positronium and its Bose-Einstein Condensation (BEC)

Positronium (Ps)

- The bound state of an electron (e-) and a positron (e+).
- The lightest and exotic atom.

Ps BEC

$$n\lambda_D^3 = n \left(\frac{2\pi\hbar^2}{mk_B T_C} \right)^{3/2} = 2.612, \quad (1)$$

n : number density,
 λ_D : de Broglie wave length,
 \hbar : reduced Planck constant,
 m : atom mass,
 k_B : Boltzmann constant,
 T_C : BEC critical temperature.

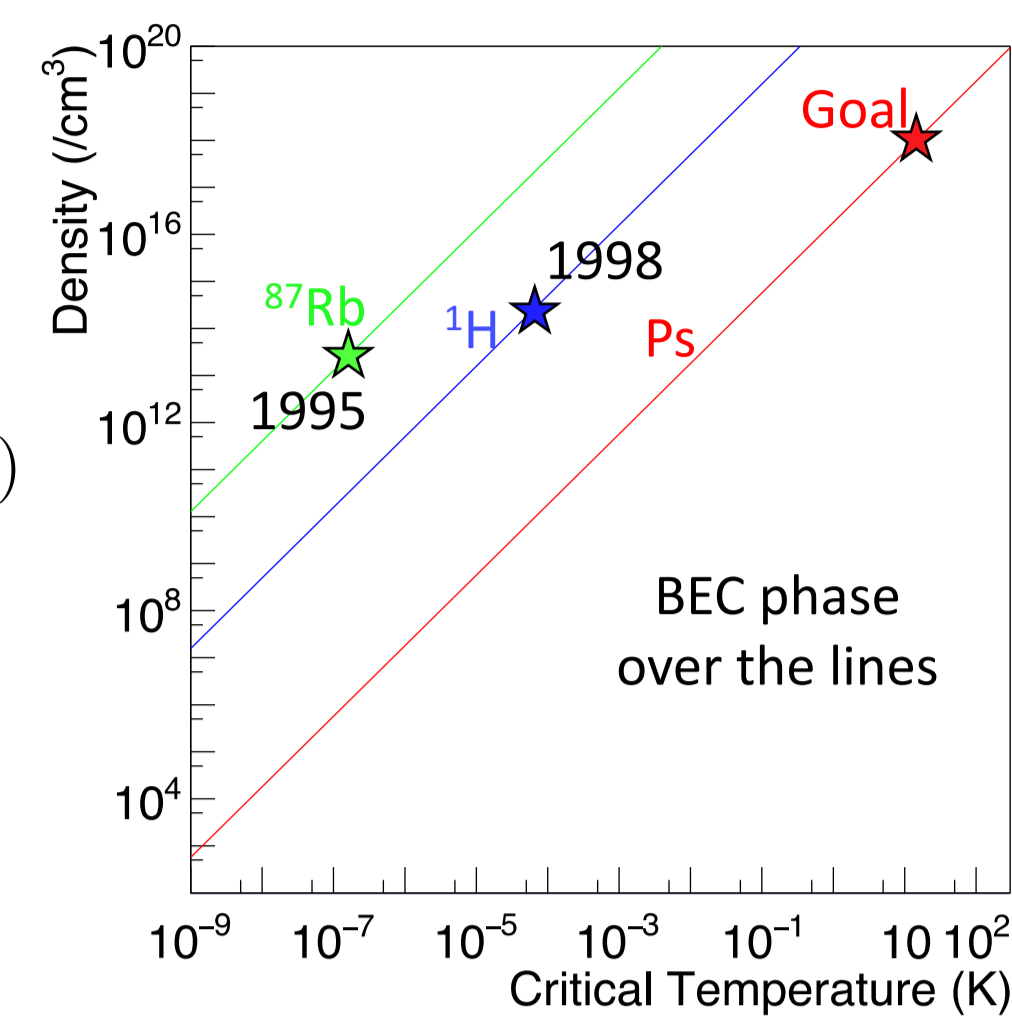


Fig. 1 BEC critical temperatures.

High critical temperature.
 Our target: **14 K @ 10¹⁸/cm³**

Scientific goal and application

- Antimatter gravity measurement
- Gamma-ray laser

Monte Carlo simulation

Setup

- 5 keV, 10⁷ e+/bunch beam.
- Silica cavity (1 K) and cooling lasers.
- (100 nm)³ cube internal void traps 4000 fully polarized Ps ($n = 4 \times 10^{18} \text{ cm}^{-3}$).

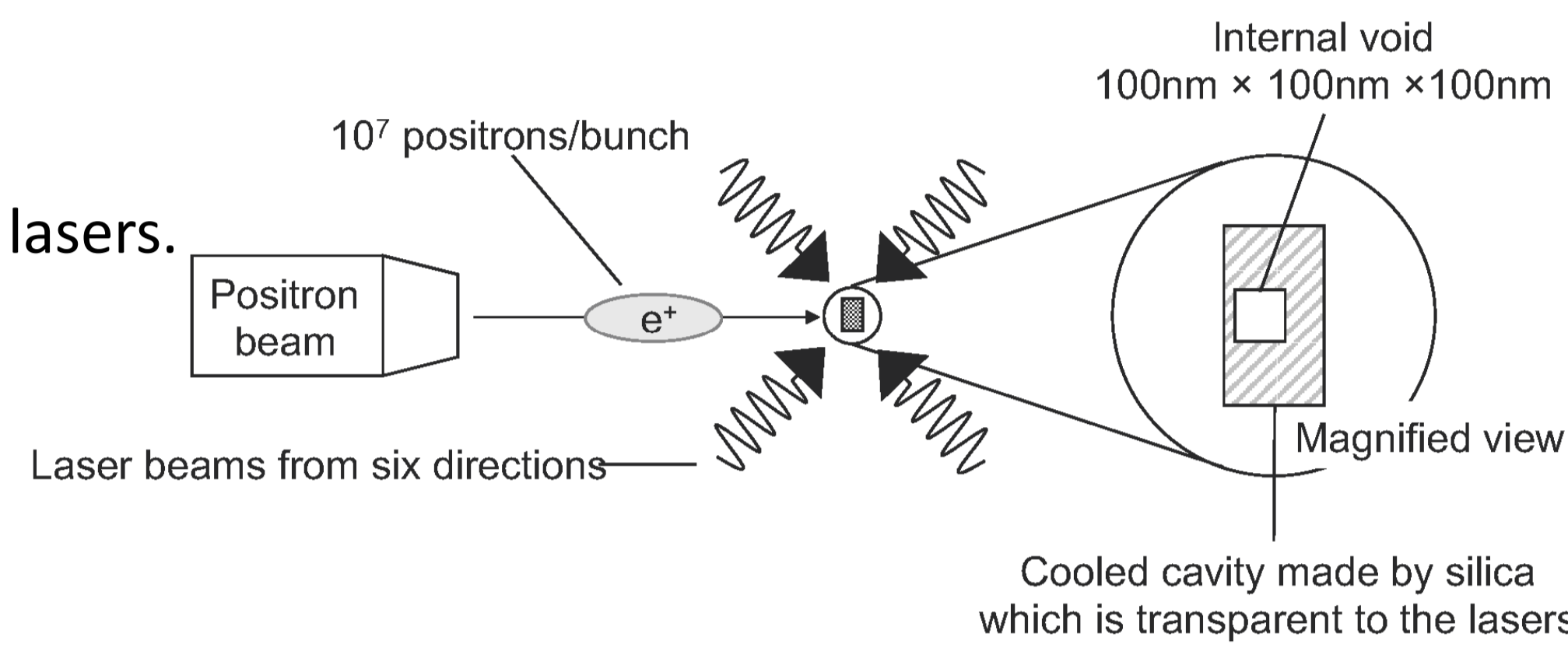


Fig. 2 Conceptual view of the setup.

Initial conditions

- 10⁴ fully polarized Ps.
- 0.8 eV monochromatic kinetic energy.
- Isotropic velocities, uniform position distributions.

Method

- Velocity and internal state of every Ps are traced at the same time (brute-force method).
- Time evolutions are calculated with short time steps.
- Random numbers are generated at each step to determine whether and which interaction each Ps does.
- Step size is only ~0.1 ps, which is short so that probabilities of all interactions are less than 1%.
- One execution without averaging.
- Positions of Ps are not traced. 100 nm cube is much smaller than the laser beam size.
- Initial Ps number of 10⁴ is larger than 4 x 10³, which is the assumed number of fully polarized Ps survived in the silica cavity, in order to decrease the statistical uncertainty. Ps-Ps scattering rate is scaled to compensate the difference.

Interactions

1. Thermalization

- Ps lose energy by collisions with silica cavity walls.
- Each Ps kinetic energy is assumed to evolve by the same way as the mean kinetic energy.
- The following differential equation (2) is solved by the Euler method.

$$\frac{dE}{dt} = -\frac{2}{LM} \sqrt{2m_{Ps}E} \left(E - \frac{3}{2}k_B T \right), \quad (2)$$

E : Ps mean kinetic energy,
 t : time since Ps formation,
 $L = \sqrt[3]{V}$ (Ps mean free path),
 V : volume of silica cavity,
 M : effective mass of silica grains,
 m_{Ps} : Ps mass,
 T : silica temperature.

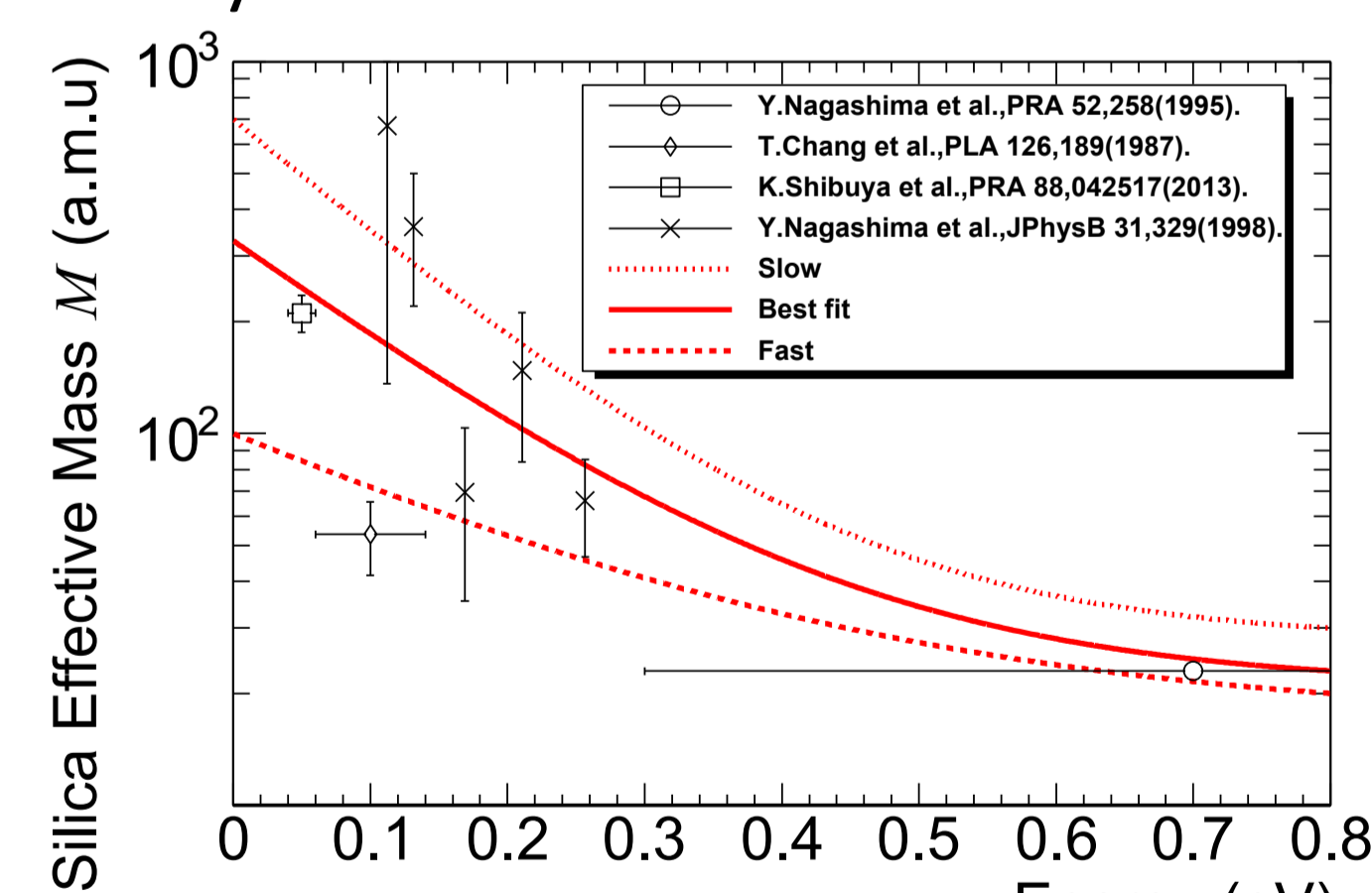


Fig. 3 Silica effective mass

- M is determined so that it represents the thermalization measurements (Fig. 3). The best-fit function is the following Eq. (3).

$$\frac{M}{\text{a.m.u.}} = 21 + 308 \exp\left(-\frac{E}{0.16 \text{ eV}}\right) \quad (3)$$

2. Ps-Ps scatterings

- Two-body elastic s -wave scatterings.
- Cross section and mean free time are calculated by Eqs. (4) and (5).

$$\sigma = 4\pi a^2, \quad (4) \quad \tau = \frac{1}{n\sigma\bar{v}}, \quad (5)$$

σ : scattering total cross section,
 $a = 0.16 \text{ nm}$ (scattering length).
 τ : scattering mean free time,
 \bar{v} : Ps mean velocity.

3. Ps decays

- 1s o-Ps annihilation (lifetime is 142 ns). Annihilated Ps are removed.
- 2p o-Ps annihilation is ignored because of long lifetime.
- Lyman- spontaneous de-excitation (3.2 ns) is included in the simulation.

4. Interaction with lasers

- Rates are calculated using Eq. (6) (right top of this poster).
- Laser intensity profile is approximated to be uniform at the peak intensity.

$$B(t, \vec{x}, \vec{v}) = \int d\omega \frac{I(t, \vec{x}, \omega)}{\hbar\omega} \cdot \frac{4}{3} \pi^2 \alpha \omega_0 |X_{12}| \cdot \frac{1}{2\pi} \times \frac{\Gamma/2}{\left[\omega \left(1 - \frac{\vec{k} \cdot \vec{v}}{c} \right) - \omega_0 \right]^2 + (\Gamma/2)^2}, \quad (6)$$

$B(t, \vec{x}, \vec{v})$: Einstein B coefficient,
 \vec{x} : Ps position,
 \vec{v} : Ps velocity,
 ω : laser frequency,
 $I(t, \vec{x}, \omega)$: laser intensity per frequency,
 α : fine structure constant,
 ω_0 : $1s - 2p$ resonant frequency,
 X_{12} : $1s - 2p$ transition matrix element,
 $\Gamma = 313 \text{ MHz}$ (natural linewidth),
 \vec{k} : laser photons' direction,
 c : speed of light.

Table 1 Assumed cooling laser specifications.

Parameters	Values
Pulse energy	40 μJ
Center frequency	1.23 PHz - $\Delta(t)$
Frequency detune $\Delta(t)$	$\Delta(0 \text{ ns}) = 300 \text{ GHz}$ $\Delta(300 \text{ ns}) = 240 \text{ GHz}$
Bandwidth (2σ)	140 GHz
Time duration (2σ)	300 ns
Beam size (2σ)	200 μm

Results

Thermalization

- Thermalization is too slow.
- Huge uncertainty because of uncertainty of M .
- Precise measurement of Ps thermalization in silica cavity is necessary.

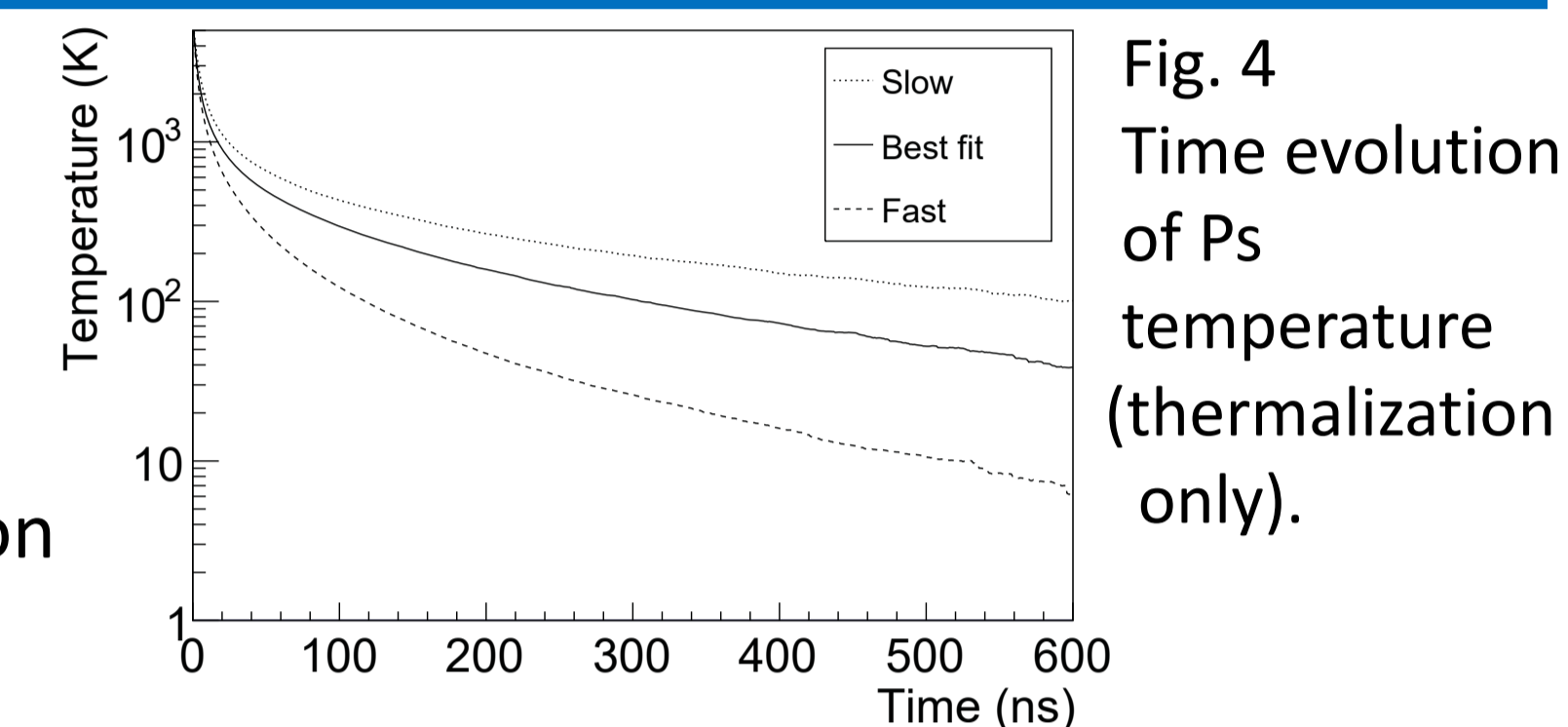


Fig. 4 Time evolution of Ps temperature (thermalization only).

Combination of Thermalization and Laser cooling

BEC can be realized by our new method (thermalization + laser cooling)!

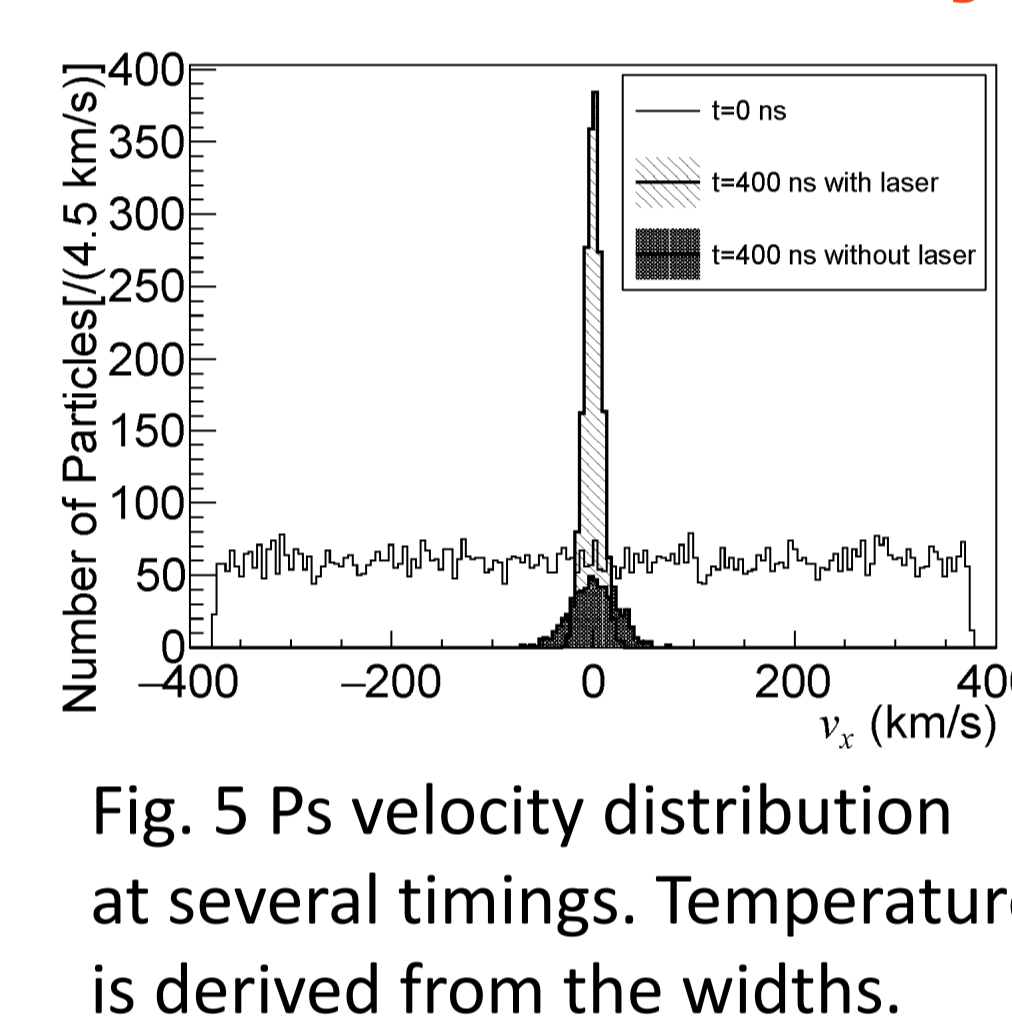


Fig. 5 Ps velocity distribution at several timings. Temperature is derived from the widths.

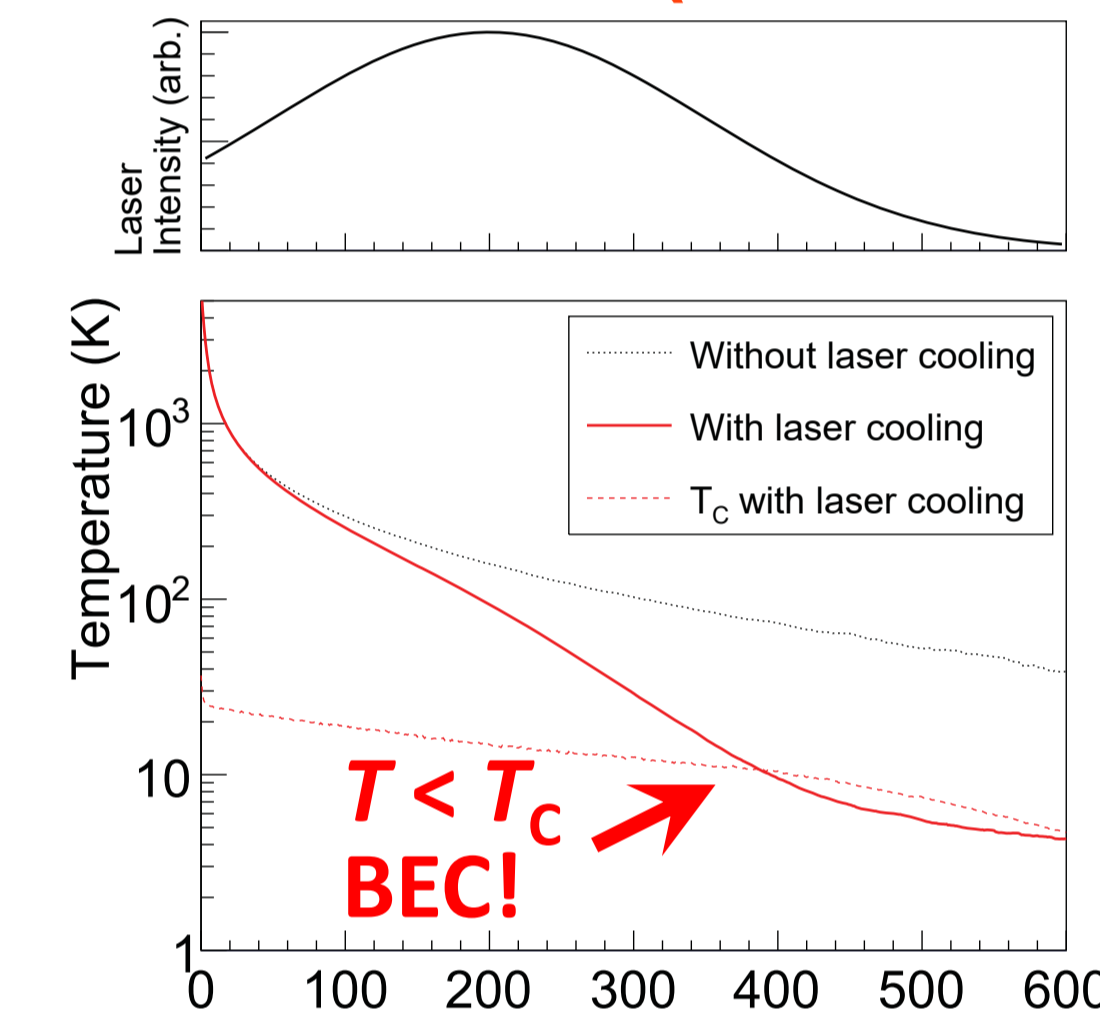


Fig. 6 Time evolution of laser intensity (top) and Ps temperature (bottom) with laser cooling (Best-fit M is used).

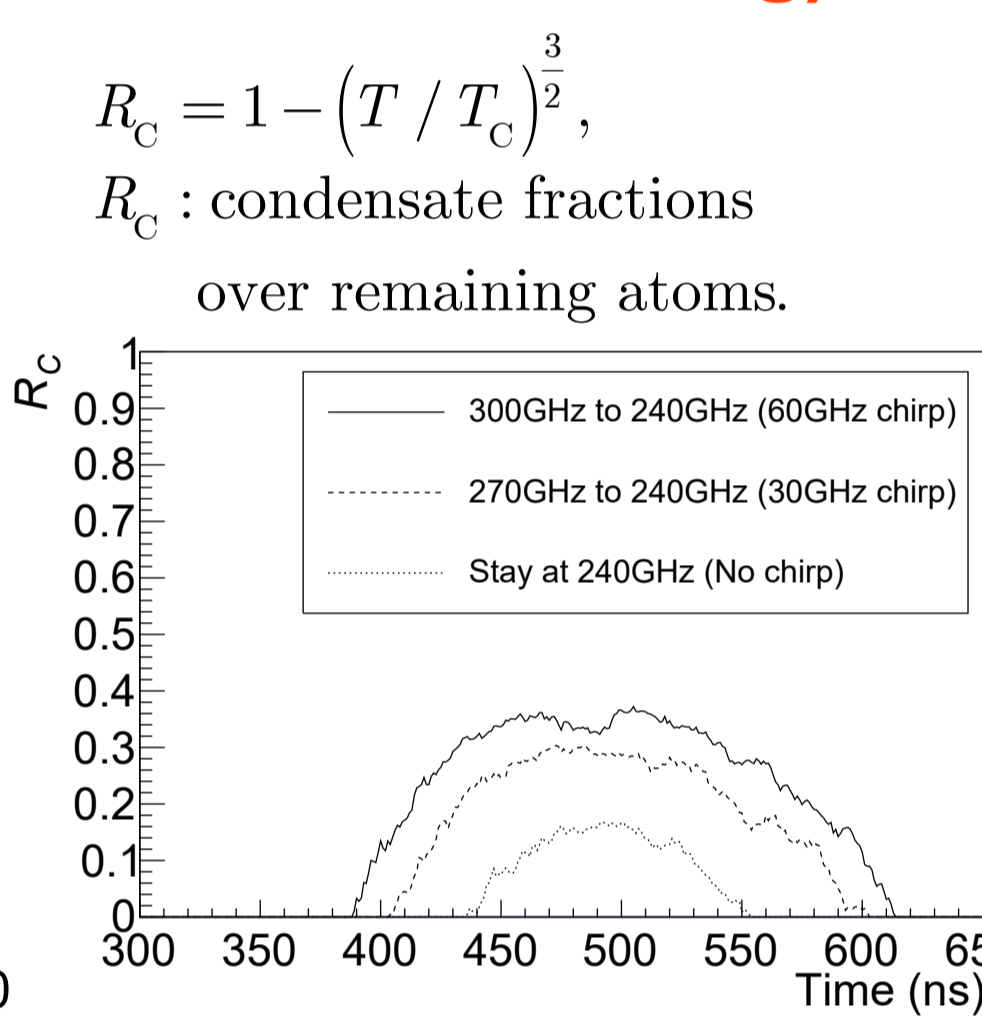
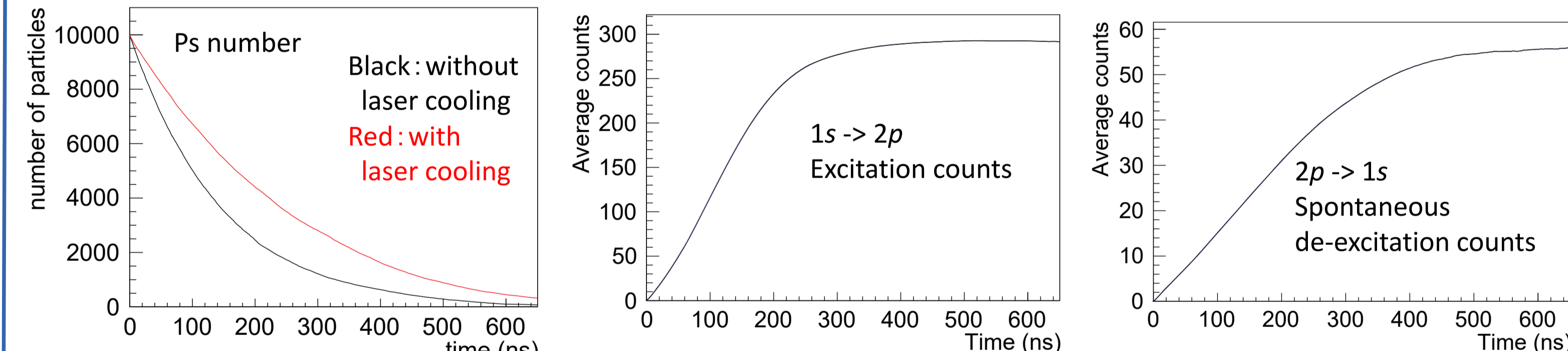


Fig. 7 Time evolution of BEC fractions with different laser frequency chirp conditions.

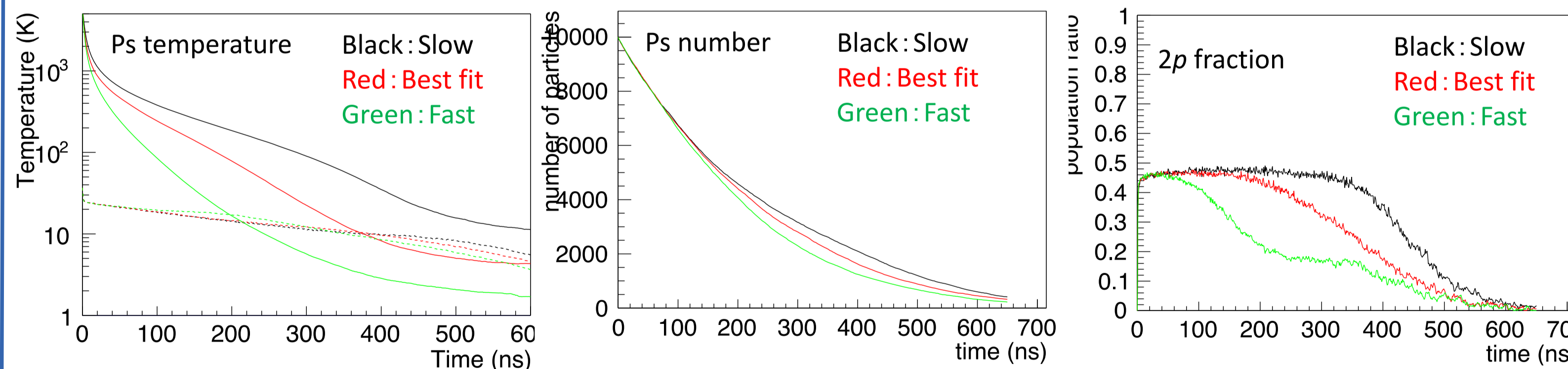
Frequency chirp is important!

Time evolutions of some parameters:

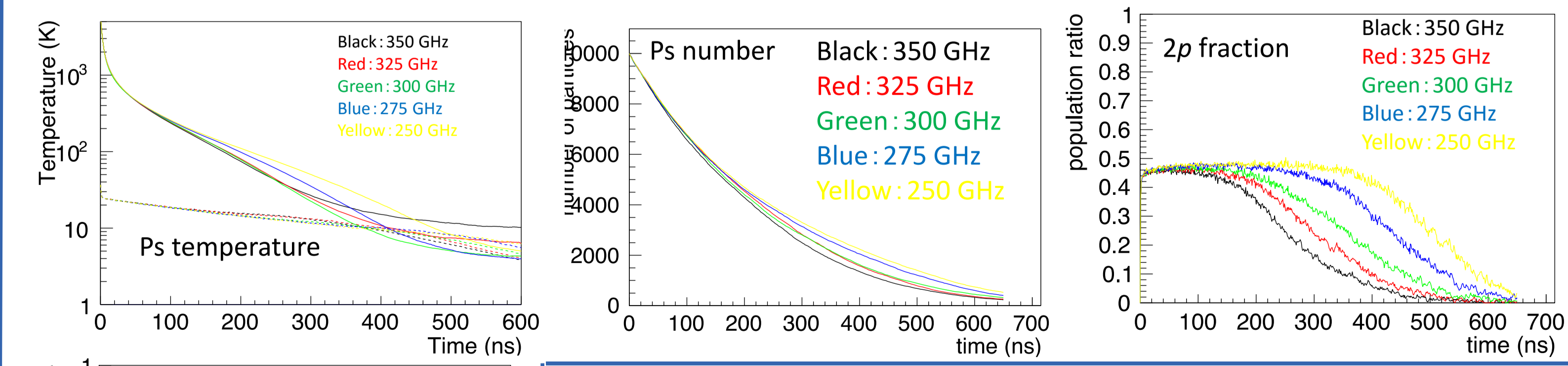


- Ps lifetime is longer with laser cooling because of 1s-2p transitions (2p has long lifetime).

Effect of uncertainty of silica effective mass M



Dependence on initial laser frequency



Summary

- We have studied a Ps cooling for BEC by a Monte Carlo simulation which includes thermalization, Ps-Ps scatterings, Ps decays and interaction with cooling lasers.
- Our simulation has demonstrated that **Ps BEC can be realized** by a **combination of the following 2 processes**, which is **our new idea** of Ps cooling:
 - (1) **Ps thermalization** in small cold silica cavity
 - (2) **Laser cooling** using 1s-2p transitions.

Reference:
 K. Shu *et al.*, "Study on cooling of positronium for Bose-Einstein condensation", J. Phys. B: At. Mol. Opt. Phys. **49**, 104001 (2016).